1 Tracy-Widom distribution

We saw that the largest eigenvalue (in magnitude) of a Hermitian Wigner matrix is around $2\sqrt{n}$. In fact, under some assumptions, we can even show that there is a central limit theorem. For simplicity we will consider only GUE.

Let Ai(x) denote the Airy function which is the unique solution of the Airy equation with prescribed asymptotics

$$\operatorname{Ai}''(x) = x\operatorname{Ai}(x), \qquad \operatorname{Ai}(x) \sum \frac{e^{-\frac{2}{3}x^{3/2}}}{2\sqrt{\pi}x^{1/4}} \text{ as } x \to +\infty.$$

Let u(x) be the unique monotone decreasing solution to the Painlevé II equation with asymptotics

$$u''(x) = 2(u(x))^3 + xu(x), \qquad u(x) \sim Ai(x) \text{ as } x \to +\infty.$$

This u also satisfies

$$u(x) \sim -\sqrt{\frac{-x}{2}}(1 + O(x^{-2}))$$
 as $x \to -\infty$.

Therefore, the function

$$F(x) := \exp\left(-\int_x^\infty (y-x)(u(y))^2 \,\mathrm{d}y\right)$$

defines a distribution function, called the Tracy-Widom distribution.

Theorem 1.1. Let $\lambda_n(M)$ be the largest eigenvalue of an $n \times n$ GUE matrix. Then for all $x \in \mathbb{R}$,

$$\lim_{n \to \infty} \mathbf{P}((\lambda_n(M) - 2\sqrt{n})n^{1/6} \le x) = F(x).$$

An interesting fact is that this distribution appears in many problems in probability, and they look unrelated to random matrices at first glance. Here we will introduce some of them.

2 Ulam's problem

Consider the symmetric group S_n and let $\pi \in S_n$. For $i_1 < i_2 < \cdots < i_k$, we say that $\pi(i_1), \ldots, \pi(i_k)$ is an increasing subsequence of π of length k if $\pi(x_1) < \cdots < \pi(i_k)$. Let $\ell_n(\pi)$ denote the maximal length of all increasing subsequence of π . An increasing subsequence of π of length $\ell_n(\pi)$ is called a longest increasing subsequence of π . For example, for the permutation

$$\pi = \begin{pmatrix} 1 & 2 & 3 & 4 & 5 \\ 5 & 1 & 3 & 2 & 4 \end{pmatrix}$$

134 and 124 are longest increasing subsequences of π . We will pick a permutation uniformly at random from S_n , how does the corresponding random variable ℓ_n behave? Surprisingly, it turns out that it behaves like the largest eigenvalue of GUE.

Theorem 2.1. For any $x \in \mathbb{R}$,

$$\lim_{n \to \infty} \mathbf{P}\left(\frac{\ell_n(\pi) - 2\sqrt{n}}{n^{1/6}} \le x\right) = F(x).$$

3 Last-passage percolation

Consider the integer lattice \mathbb{Z}^2 . Let $(t_x)_{x\in\mathbb{Z}^2}$ be a family of i.i.d. random variables. An oriented path γ with vertices x_0, \ldots, x_n has the property that all coordinates of x_i are no larger than those of x_{i+1} . For an oriented path γ , we define the passage time $T(\gamma) = \sum_{x\in\gamma\setminus\{x_0\}} t_x$. Finally, we define the last-passage time between x and y by

$$T(x, y) = \max\{T(\gamma) : \gamma \text{ is an oriented path from } x \text{ to } y\}.$$

Of course, this only makes sense when there is an oriented path from x to y. Again, this model is related to Tracy-Widom distribution.

Theorem 3.1. Suppose that t_x is either exponentially distributed with parameter 1 or geometrically distributed with parameter p. Then for all $x \in \mathbb{R}$,

$$\lim_{m,n\to\infty,m/n\to\alpha\in(0,1)} \mathbf{P}\left(\frac{T((0,0),(m,n)) - g(m,n)}{\sigma(m,n)} \le x\right) = F(x),$$

where

$$g(m,n) = \begin{cases} (\sqrt{m} + \sqrt{n})^2 & \text{if } t_x \text{ is exponential}, \\ \frac{2\sqrt{mnp} + (m+n)p}{1-p} & \text{if } t_x \text{ is geometric}, \end{cases}$$

and

$$\sigma(m,n) = \begin{cases} \frac{(\sqrt{m} + \sqrt{n})^{4/3}}{(mn)^{1/6}} & \text{if } t_x \text{ is exponential,} \\ \\ \frac{p^{1/6}(\sqrt{m} + \sqrt{np})^{2/3}(\sqrt{n} + \sqrt{mp})^{2/3}}{(1-p)(mn)^{1/6}} & \text{if } t_x \text{ is geometric.} \end{cases}$$

3.1 TASEP

There are some models that are equivalent to the last-passage percolation. We will describe one of them, TASEP.

The total asymmetric simple exclusion process (TASEP) is an important stochastic interacting particle systems. At any time, each site $j \in \mathbb{Z}$ is either occupied by a particle or is empty. Let $\eta(t) = {\eta_j(t)}_{j \in \mathbb{Z}}$ be defined by $\eta_j(t) = 1$ if site j has a particle at time t and 0 otherwise. Particles can jump to the site on the right if the site on the right is empty. When the right site is empty, a jump is performed after an exponential waiting time with mean 1. All jumps are independent of each other.

Suppose that at each site $j \in \mathbb{Z}$ with $j \leq 0$, there sits a particle at time 0. The particle at site 0 is allowed to move without any restriction, but other particles may sometimes be blocked.

TASEP is related to the last-passage percolation for the following reason. Consider lastpassage percolation with exponential weights. The procession of the first particle in TASEP is the same as last-passage time from (0,0) along the x-direction. For the second level (that is, y = 1), the number of sites reached in last-passage percolation corresponds to the number of steps the second particle moved in TASEP. In general, the *n*-th step of the *k*-th particle in TASEP corresponds to the site (n - 1, k) being reached from (0, 0). This shows why these two models are equivalent. In fact, the formula of g (the exponential case) in Theorem 3.1 was actually obtained using TASEP.

References

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